QUANTUM MECHANICS Lecture 17

Enrico Iacopini

## QUANTUM MECHANICS Lecture 17 Linear Operators in $\mathcal{H}$ and Physical Observables

Enrico Iacopini

November 6, 2019

D. J. Griffiths: paragraphs 3.2

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In the previous lecture, we have seen that, if  $\{e_j\}$  and  $\{f_k\}$  are two orthonormal bases of a Hilbert space, then there exist two square matrices A and B such that<sup>1</sup>

$$\mathbf{f}_{\mathbf{k}} = A_{jk} \, \mathbf{e}_{\mathbf{j}}; \quad \mathbf{e}_{\mathbf{j}} = B_{kj} \, \mathbf{f}_{\mathbf{k}}$$
  
with  $A B = I; \quad (\Rightarrow A = B^{-1}; \quad B = A^{-1}).$ 

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<sup>&</sup>lt;sup>1</sup>We are using the Einstein convention for which, when an index appears twice, it implies summation of that index over all its possible values.

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- The matrices A and B describe also the effect of the basis change on the components of a generic vector.
- In fact, if the components of a generic vector ν in the basis {e<sub>j</sub>} are {α<sub>j</sub>} and in the basis {f<sub>k</sub>} are {β<sub>k</sub>}, then

$$\mathbf{v} \equiv \alpha_j \mathbf{e}_{\mathbf{j}} = \alpha_j B_{kj} \mathbf{f}_{\mathbf{k}} \equiv \beta_k \mathbf{f}_{\mathbf{k}} \implies$$
$$\Rightarrow \beta_k = B_{kj} \alpha_j$$

$$\mathbf{v} \equiv \beta_k \mathbf{f}_{\mathbf{k}} = \beta_k A_{jk} \mathbf{e}_{\mathbf{j}} \equiv \alpha_j \mathbf{e}_{\mathbf{j}} \Rightarrow$$
$$\Rightarrow \alpha_j = A_{jk} \beta_k$$

where we have used the fact that, in a basis, the components of a vector are unique.

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It can be shown that the linear space of the square-integrable functions  $\psi(x)$  is a Hilbert space with the following scalar product

$$<\psi_1|\psi_2>\equiv\int dx\;\psi_1(x)^*\,\psi_2(x)$$

Clearly, the dimension of H is not finite, which means that the cardinality of any basis has to be numerable. QUANTUM MECHANICS Lecture 17

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According to what we have seen, if  $\{\psi_j\}$  is an orthonormal basis of the Hilbert space  $\mathcal{H}$  of the square integrable functions in one real variable, then

$$<\psi_i|\psi_j>\equiv\int dx\;\psi_i(x)^*\,\psi_j(x)=\delta_{ij}$$

2 if  $\psi \in \mathcal{H}$ 

1

$$\psi = \sum\limits_j < \psi_j | \psi > \ \psi_j \equiv c_j \ \psi_j$$

where

$$c_j = \int dx \, \psi_j(x)^* \, \psi(x)$$

) and, if  $\psi$  is normalized

 $\sum |c_j|^2 = 1 \quad \text{and} \quad z \to z \to z$ 

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- So much for the moment about the physical states and their representation as vectors of a Hilbert space ...
- If the appropriate mathematical structure in which to describe the states of a physical system in QM is the Hilbert space, which is the corresponding description of the physical observables ?

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- If the appropriate mathematical structure in which to describe the states of a physical system in QM is the Hilbert space, which is the corresponding description of the physical observables ?

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We have already seen that, to evaluate, for instance, the expectation value of a physical observable Q = Q(x, p) on a physical state described by the w.f.  $\Psi(x, t)$ , we need to calculate the integral

$$\int dx \, \Psi(x,t)^* \left[ \widehat{Q}\left(x,-i\hbarrac{\partial}{\partial x}
ight) \, \Psi(x,t) 
ight]$$

in which we can now recognize the scalar product of  $\Psi$  and  $\hat{Q}\Psi$ , being  $\hat{Q}$  a linear operator.

But, which is the proper definition of a linear operator in a Hilbert space ? QUANTUM MECHANICS Lecture 17

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- In general, a linear operator Q acting on the vectors of a Hilbert space must be such that
  - $\forall a \in \mathcal{H}, \ \hat{Q}a \in \mathcal{H}$
  - $\hat{Q}(\mathbf{a} + \mathbf{b}) = \hat{Q}\mathbf{a} + \hat{Q}\mathbf{b}$
  - $\hat{Q}(\alpha \mathbf{a}) = \alpha \, \hat{Q} \mathbf{a}$
- We have seen that, once we have choosen an orthonormal basis {e<sub>i</sub>}, a vector is completely determined by its numerical (complex) components.

Do we have something similar for a linear operator ? QUANTUM MECHANICS Lecture 17

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Do we have something similar for a linear operator ?

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Let us consider an operator  $\hat{Q}$  and a generic vector  $\mathbf{a} = \alpha_i \mathbf{e}_i$ , where  $\{\mathbf{e}_i\}$  is a suitable orthonormal basis. We have

$$\hat{Q}\mathbf{a} = \hat{Q}(\alpha_i \mathbf{e}_i) = \alpha_i \hat{Q}\mathbf{e}_i$$

But, for any vector  $\mathbf{e}_i$  of the basis, we will have

$$\hat{Q}\mathbf{e}_i = Q_{ji}\mathbf{e}_j$$
 where  $Q_{ji} = \langle \mathbf{e}_j | \hat{Q}\mathbf{e}_i \rangle \in C$ 

and therefore

$$\hat{Q}\mathbf{a} = \alpha_i \,\hat{Q}\mathbf{e}_i = \alpha_i \,Q_{ji}\mathbf{e}_j = (Q_{ji}\alpha_i)\mathbf{e}_j \implies \Rightarrow <\mathbf{e}_j |\hat{Q}\mathbf{a}\rangle = Q_{ji}\alpha_i = Q_{ji} < \mathbf{e}_i |\mathbf{a}\rangle$$

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In conclusion, in a given basis  $\{\mathbf{e}_i\}$ , the action of a linear operator  $\hat{Q}$  is described by a suitable complex matrix  $Q_{ji} = \langle \mathbf{e}_j | \hat{Q} \mathbf{e}_i \rangle$ , such that the components of a generic vector  $\hat{Q}\mathbf{a}$  are given by

$$<\mathbf{e}_{j}|\hat{Q}\mathbf{a}>=Q_{ji}lpha_{i}$$

where the  $\{\alpha_i\}$  are the component of the vector **a**, in the same basis.

Let us see, now, what happens to the matrix Q<sub>ji</sub>, representing a given linear operator Q̂, when we change the vector basis.

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Let us see, now, what happens to the matrix Q<sub>ji</sub>, representing a given linear operator Q̂, when we change the vector basis.

• We have already seen that, if  $\{\mathbf{e}_j\}$  and  $\{\mathbf{f}_k\}$  be two orthonormal bases, and

$$\mathbf{f}_k = A_{jk} \mathbf{e}_j; \quad \mathbf{e}_j = A_{kj}^{-1} \mathbf{f}_k$$

then, if  $\mathbf{a}$  is a generic vector such that

$$\mathbf{a} = \alpha_j \, \mathbf{e}_j = \beta_k \, \mathbf{f}_k$$

we have

$$lpha_j = A_{jk} \, eta_k; \quad eta_k = A_{kj}^{-1} \, lpha_j$$

2 Let

$$<\mathbf{e}_{j}|\hat{Q}\mathbf{a}>\equiv Q_{ji}\alpha_{i}$$
  
 $<\mathbf{f}_{k}|\hat{Q}\mathbf{a}>\equiv \tilde{Q}_{kt}\beta_{t}$ 

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$$\mathbf{f}_k = A_{jk} \mathbf{e}_j; \quad \mathbf{e}_j = A_{kj}^{-1} \mathbf{f}_k$$

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we have

$$\alpha_j = A_{jk} \beta_k; \quad \beta_k = A_{kj}^{-1} \alpha_j$$

2 Let

$$\langle \mathbf{e}_{j} | \hat{Q} \mathbf{a} \rangle \equiv Q_{ji} \alpha_{i}$$
  
 $\langle \mathbf{f}_{k} | \hat{Q} \mathbf{a} \rangle \equiv \tilde{Q}_{kt} \beta_{t}$ 

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Therefore

$$<\mathbf{f}_{k}|\hat{Q}\mathbf{a}>=\tilde{Q}_{kt}\beta_{t}=\tilde{Q}_{kt}A_{tj}^{-1}\alpha_{j}\equiv\left(\tilde{Q}A^{-1}
ight)_{kj}\alpha_{j}$$

$$\langle \mathbf{f}_k | \hat{Q} \mathbf{a} \rangle = \langle A_{ik} \mathbf{e}_i | \hat{Q} \mathbf{a} \rangle = A_{ik}^* \langle \mathbf{e}_i | \hat{Q} \mathbf{a} \rangle =$$

$$= A_{ik}^* Q_{ij} \alpha_j = A_{ki}^+ Q_{ij} \alpha_j = (A^+ Q)_{kj} \alpha_j$$

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Therefore

$$<\mathbf{f}_{k}|\hat{Q}\mathbf{a}>=\tilde{Q}_{kt}eta_{t}= ilde{Q}_{kt}A_{tj}^{-1}lpha_{j}\equiv\left( ilde{Q}A^{-1}
ight)_{kj}lpha_{j}$$

On the other hand  $\langle \mathbf{f}_k | \hat{Q} \mathbf{a} \rangle = \langle A_{ik} \mathbf{e}_i | \hat{Q} \mathbf{a} \rangle = A_{ik}^* \langle \mathbf{e}_i | \hat{Q} \mathbf{a} \rangle =$  $= A_{ik}^* Q_{ij} \alpha_j = A_{ki}^+ Q_{ij} \alpha_j = \left(A^+ Q\right)_{kj} \alpha_j$ 

where the matrix  $A^+ \equiv (A^t)^*$  is the transposed complex conjugate of the matrix A, called **the hermitian conjugate** of A.

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Therefore

$$<\mathbf{f}_{k}|\hat{Q}\mathbf{a}>= ilde{Q}_{kt}eta_{t}= ilde{Q}_{kt}A_{tj}^{-1}lpha_{j}\equiv\left( ilde{Q}A^{-1}
ight)_{kj}lpha_{j}$$

On the other hand

$$<\mathbf{f}_{k}|\hat{Q}\mathbf{a}> =  =A_{ik}^{*}<\mathbf{e}_{i}|\hat{Q}\mathbf{a}> =$$
$$=A_{ik}^{*}Q_{ij}\alpha_{j}=A_{ki}^{+}Q_{ij}\alpha_{j}=\left(A^{+}Q\right)_{kj}\alpha_{j}$$

where the matrix A<sup>+</sup> ≡ (A<sup>t</sup>)\* is the transposed complex conjugate of the matrix A, called *the hermitian conjugate* of A.
Given of the arbitrarity of the α<sub>j</sub>, we can conclude that it must be

$$\tilde{Q} A^{-1} = A^+ Q \quad \Leftrightarrow \quad \tilde{Q} = A^+ Q A$$

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